

## Article

**On the Horndeski's Formula for the Lanczos Invariant**H. N. Núñez-Yépez<sup>1</sup>, A. L. Salas-Brito<sup>2</sup> & J. López-Bonilla<sup>3\*</sup><sup>1</sup>Departamento de Física, UAM-I, Apdo. Postal 55-534, Iztapalapa 09340, CDMX, México<sup>2</sup>Lab. de Sistemas Dinámicos, Depto. de Ciencias Básicas, UAM-A, Apdo. Postal 21-267, Coyoacán CP 0400, CDMX, México<sup>3</sup>ESIME-Zacatenco, IPN, Edif. 5, 1er. Piso, Lindavista 07738, CDMX, México**Abstract**

Horndeski wrote the Lagrangian  $\sqrt{-g} {}^*R^{\mu\nu\alpha\beta} R_{\mu\nu\alpha\beta}$  as an exact divergence; here we exhibit that this formula is useful to study the existence of non-null constant vectors in  $R_4$ .

**Keywords:** Lanczos scalar, Horndeski's formula, Lanczos potential.

**1. Introduction**

Here we consider the Lanczos invariant [1]:

$$K_2 = {}^*R^{\mu\nu\alpha\beta} R_{\mu\nu\alpha\beta}, \quad (1)$$

which is the contraction of the Riemann tensor with its double dual [2]. If we employ the Lagrangian  $L = \sqrt{-g} K_2$  in the Hilbert type variational principle  $\delta \int L d^4x = 0$ , we obtain  $0 = 0$ , hence one suspects that the density  $L$  is an exact divergence for any  $R_4$ :

$$\sqrt{-g} K_2 = (\sqrt{-g} C^\mu)_{,\mu} \quad (2)$$

where  $_{,\mu} = \frac{\partial}{\partial x^\mu}$ . This suspicion turned out to be correct because Buchdal [5, 6] and Goenner-Kohler [7] got non-tensorial expressions for  $C^\nu$ . From (2) Lanczos [8] proved that the Weyl tensor is generated by a potential of third order.

In Sec. 2 we indicate the Horndeski's tensorial expression [9] for  $C^\alpha$  and we exhibit that it is useful to analyze the presence of non-null constant vectors in a given space-time.

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## 2. Horndeski's formula for $\sqrt{-g} K_2$

It is interesting to note that the vector field  $C^\mu$  obtained by Horndeski [9] contains an arbitrary non-null vector  $A^\nu$ , that without loss of generality we take unitary and time-like, that is,  $A^\alpha A_\alpha = 1$ , then:

$$C^\mu = 8 \left( {}^*R^{*\mu\nu}{}_{\alpha\beta} + \frac{1}{3} \delta_{\tau\beta\varphi\alpha}^{\lambda\nu\theta\mu} A^\tau{}_{;\lambda} A^\varphi{}_{;\theta} \right) A^\beta{}_{;\tau} A^\alpha, \quad (3)$$

verifying (2), with the participation of the generalized Kronecker delta [4]. We observe that  $(\sqrt{-g} C^\nu)_{;\nu} = (\sqrt{-g} C^\nu)_{;\nu}$  because (3) has tensorial character.

Therefore, when  $K_2 \neq 0$  the space-time does not accept non-null constant vectors, in other words, the existence of a non-null vector such that  $A^\mu{}_{;\nu} = 0$  implies  $K_2 = 0$  via (2) and (3) (which occurs, for example, in the Gödel solution [2, 10, 11]). The metrics of Schwarzschild, Taub, C, Kerr, .... have  $K_2 \neq 0$  [2], hence these space-times do not admit non-null constant vectors.

Thus, the Lagrangian  $\sqrt{-g} K_2$  is an ordinary divergence, however, it can contribute [12, 13] to the gravitational energy-momentum distribution. For empty spaces, in [14] there is a tensor formula for  $C^\mu$  in terms of the Lanczos potential [8].

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