

A Potential for the Lanczos Spinor in Arbitrary Spacetimes of Petrov Types III, N, and O

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Abstract

Andersson-Edgar proved that the Lanczos spinor can be generated via $L_{ABCD} = \nabla^E_{\dot{D}} T_{ABCE}$. Here we determine the potential T_{ABCE} for any R_4 of Petrov types O, N, and III.

Keywords: Lanczos spinor, spin coefficients.

1. Introduction

In [1] was obtained the Lanczos generator [2, 3] for an arbitrary spacetime of Petrov type N, with the following Newman-Penrose components:

$$\begin{aligned} \Omega_0 &= \frac{\kappa}{2}, & \Omega_3 &= -\frac{\lambda}{2}, & \Omega_4 &= \frac{\sigma}{2}, & \Omega_7 &= -\frac{\nu}{2}, \\ \Omega_1 &= \frac{\rho}{6}, & \Omega_2 &= -\frac{\pi}{6}, & \Omega_5 &= \frac{\tau}{6}, & \Omega_6 &= -\frac{\mu}{6}, \end{aligned} \tag{1}$$

in terms of the spin coefficients [4, 5] associated to the canonical null tetrad [6]:

$$l^\mu \leftrightarrow o^A o^{\dot{B}}, \quad n^\mu \leftrightarrow l^A l^{\dot{B}}, \quad m^\mu \leftrightarrow o^A l^{\dot{B}}, \quad \bar{m}^\mu \leftrightarrow l^A o^{\dot{B}}; \tag{2}$$

hence for the type N the Lanczos spinor [7, 8] is given by:

$$\begin{aligned} L_{ABCD} &= \frac{1}{6} [3 o_A o_B o_C (-\nu o_{\dot{D}} + \lambda l_{\dot{D}}) + (o_A o_B l_C + (o_A * l_B) o_C) (\mu o_{\dot{D}} - \pi l_{\dot{D}}) + \\ &+ (l_A l_B o_C + (o_A * l_B) l_C) (\tau o_{\dot{D}} - \rho l_{\dot{D}}) + 3 l_A l_B l_C (-\sigma o_{\dot{D}} + \kappa l_{\dot{D}})]. \end{aligned} \tag{3}$$

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On the other hand, Andersson-Edgar [9-11] proved that any Lanczos spinor can be generated via the relation:

$$L_{ABCD} = \nabla^E_D T_{ABCE}, \quad T_{ABCE} = T_{(ABC)E}. \quad (4)$$

In Sec. 2 we give an explicit expression for the potential T_{ABCE} in terms of the dyad (o^F, ι^G) , for the Lanczos spinor (3).

Let's remember that [1]:

$$(\Omega_r)_{Type III} = 2 (\Omega_r)_{Type N}, \quad r = 0, \dots, 7, \quad (5)$$

in the canonical tetrad for the type III, therefore:

$$(T_{ABCE})_{Type III} = 2 (T_{ABCE})_{Type N}. \quad (6)$$

2. Potential for the Lanczos spinor

In the spin coefficients formalism [4] are very known the following relations:

$$\begin{aligned} \nabla^A_{\dot{B}} o_A &= (\tau - \beta) o_{\dot{B}} + (\varepsilon - \rho) \iota_{\dot{B}}, & \nabla^A_{\dot{B}} \iota_A &= (\gamma - \mu) o_{\dot{B}} + (\pi - \alpha) \iota_{\dot{B}}, \\ o_A \nabla^A_{\dot{B}} o_C &= (-\beta o_{\dot{B}} + \varepsilon \iota_{\dot{B}}) o_C + (\sigma o_{\dot{B}} - \kappa \iota_{\dot{B}}) \iota_C, \\ o_A \nabla^A_{\dot{b}} \iota_C &= (-\mu o_{\dot{B}} + \pi \iota_{\dot{B}}) o_C + (\beta o_{\dot{B}} - \varepsilon \iota_{\dot{B}}) \iota_C, \\ \iota_A \nabla^A_{\dot{B}} o_C &= (-\gamma o_{\dot{B}} + \alpha \iota_{\dot{B}}) o_C + (\tau o_{\dot{B}} - \rho \iota_{\dot{B}}) \iota_C, \\ \iota_A \nabla^A_{\dot{B}} \iota_C &= (-\nu o_{\dot{B}} + \lambda \iota_{\dot{B}}) o_C + (\gamma o_{\dot{B}} - \alpha \iota_{\dot{B}}) \iota_C, \end{aligned} \quad (7)$$

hence:

$$\begin{aligned} \nabla^E_{\dot{D}} [(o_A o_B \iota_C + (o_A * \iota_B) o_C) \iota_E] &= 3 o_A o_B o_C (-\nu o_{\dot{D}} + \lambda \iota_{\dot{D}}) + \\ &+ (o_A o_B \iota_C + (o_A * \iota_B) o_C) (-\mu o_{\dot{D}} + \pi \iota_{\dot{D}}) + 2 (\iota_A \iota_B o_C + (o_A * \iota_B) \iota_C) (\tau o_{\dot{D}} - \rho \iota_{\dot{D}}), \\ \nabla^E_{\dot{D}} [(\iota_A \iota_B o_C + (o_A * \iota_B) \iota_C) o_E] &= 3 \iota_A \iota_B \iota_C (\sigma o_{\dot{D}} - \kappa \iota_{\dot{D}}) + \\ &+ (\iota_A \iota_B o_C + (o_A * \iota_B) \iota_C) (\tau o_{\dot{D}} - \rho \iota_{\dot{D}}) + 2 (o_A o_B \iota_C + (o_A * \iota_B) o_C) (-\mu o_{\dot{D}} + \pi \iota_{\dot{D}}), \end{aligned} \quad (8)$$

then it is immediate that (4) implies (3) for:

$$T_{ABCE} = \frac{1}{6} [(o_A o_B l_C + (o_A * l_B) o_C) l_E - (l_A l_B o_C + (o_A * l_B) l_C) o_E], \quad (9)$$

where it is evident the symmetry (4) in ABC .

The result (9) is applicable to the type III in according with (6). The type O accepts as Lanczos spinor to any multiple of (3) [1], thus (9) also is valid for this type.

Received November 23, 2016; Accepted December 25, 2016

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